

Rethinking the QCD collisional energy loss

ANDRÉ PESHIER

Institute for Theoretical Physics, Giessen University, Germany

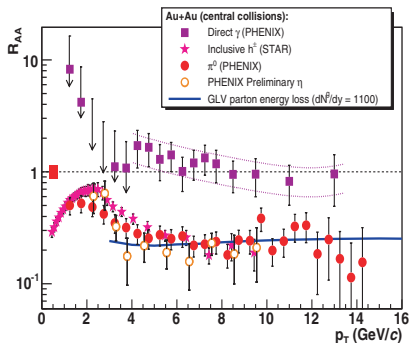
– Hot Quarks '06, Villasimius (Italy), May 2006 –

- 1 Bjorken's formula
- 2 What is the relevant coupling?
- 3 dE_{col}^{QCD} / dx

Why could it be interesting?

piece of evidence for **QGP at RHIC**

- nuclear modification factor, hadrons vs. photons



⇒ jet quenching ⇒ **energy loss**

- final state interactions
- radiative loss
[BDMPS], [GLV], [W], [SW], ...

| expected | 'seen' |
|--|---|
| $\hat{q} \sim 2\text{GeV}^2/\text{fm}$ | $\hat{q} \sim 10\text{GeV}^2/\text{fm}$ |
| $\frac{dN_{qg}}{dy} \sim 1000$ | $\frac{dN_{qg}}{dy} \sim 2000$ |

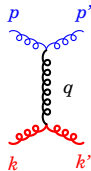
'there must be something else ...'

- **collisional loss**
[Bjorken], [Braaten, Thoma], ...
early estimates $\sim 0.3\text{GeV}/\text{fm}$
now few people say it could be larger

mean energy loss of jet j with energy E in (thermal) medium of scatterers s

$$\frac{dE_j}{dx} = \sum_s \int_{k^3} \rho_s(k) \Phi \int dt \frac{d\sigma_{js}}{dt} \Delta E, \quad \Delta E = E - E'$$

- t -channel dominates $\frac{d\sigma_{js}}{dt} = 2\pi C_{js} \frac{\alpha^2}{t^2}$



- hard jets ($E, E' \gg k \sim T$): $-t = q^2 = 2(1 - \cos\theta)k\Delta E$
 $\Phi = 1 - \cos\theta$

Bjorken evaluates

$$\Phi \int dt \frac{d\sigma_{js}}{dt} \Delta E = \frac{2\pi C_{js} \alpha^2}{-2k} \int_{t_1}^{t_2} \frac{dt}{t} = \frac{\pi C_{js} \alpha^2}{k} \ln \frac{t_1}{t_2}$$

cut-offs in t -integral ... 'physically intuitive' ...

- soft: exchanged gluon is Debye screened, $-t_2 = \mu^2 \sim m_D^2 \sim \alpha T^2$
- hard: kinematic constraint $\Delta E_{max} \sim E/2$, $-t_1 = (1 - \cos\theta)kE$

$$\Rightarrow \frac{dE_j^B}{dx} = \pi\alpha^2 \sum_s C_{js} \int_{k^3} \frac{\rho_s}{k} \ln \frac{(1 - \cos\theta)kE}{\mu^2}$$

Bjorken 'calculates' $\ln(\%) \rightarrow \ln(2\langle k \rangle E / \mu^2)$ with $\langle k \rangle \rightarrow 2T$, arriving at

$$\frac{dE_{q,g}^B}{dx} = \left(\frac{2}{3}\right)^{\pm 1} 2\pi \left(1 + \frac{1}{6} n_f\right) T^2 \alpha^2 \ln \frac{4TE}{\mu^2}$$

[Bjorken '82]

relativistic adaption of (QED) Bethe-Bloch formula

NB: HTL calculations [Thoma et al.] to fix cut-offs: $\mu^* = m_D$

What is α ?

Loop corrections (massless QED)

beyond tree level



1 $T = 0$

eff. matrix element $\mathcal{M} = \frac{\alpha}{P^2 - \Pi_{\text{vac}}}$

bare coupling

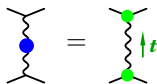
unrenormalized self-energy

$$\Pi_{\text{vac}}(P^2) = 4\pi\beta_0\alpha \left[\epsilon^{-1} - \ln \frac{-P^2}{\bar{\mu}^2} \right] P^2$$

coupling renormalization

- **measure at $P^2 = t_e$:** $\mathcal{M}(t_e) = \frac{1}{t_e} \frac{\alpha}{1 - 4\pi\beta_0\alpha[\epsilon^{-1} - \ln(-t_e/\bar{\mu}^2)]} \equiv \frac{\alpha(t_e)}{t_e}$ observable!
- $\Rightarrow \alpha^{-1}(t_e) = \alpha^{-1} - 4\pi\beta_0 [\epsilon^{-1} - \ln(-t_e/\bar{\mu}^2)]$ ‘running coupling’
- **predict at t :** $\alpha^{-1}(t) = \alpha^{-1}(t_e) + 4\pi\beta_0 \ln(t/t_e)$ (RG), or

$$\alpha(t) = [4\pi\beta_0 \ln(|t|/\Lambda^2)]^{-1}$$



Loop corrections (massless QED)

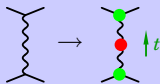
2 $T > 0$

additional contribution to selfenergy components ($i = t, l$)

$$\Pi^i(p_0, p) = 4\pi\beta_0\alpha \left[(\epsilon^{-1} - \ln(-P^2/\bar{\mu}^2)) P^2 + f^i(p_0, p) \right]$$

rewrite ('renormalize') in-medium scattering matrix

$$\begin{aligned} \frac{\alpha}{P^2 - \Pi^i} &= \frac{P^{-2}}{\alpha^{-1} - 4\pi\beta_0[\epsilon^{-1} - \ln(-P^2/\bar{\mu}^2) + f^i/P^2]} \\ &= \frac{P^{-2}}{\alpha^{-1}(P^2) - 4\pi\beta_0 f^i/P^2} = \frac{\alpha(P^2)}{P^2 - 4\pi\beta_0\alpha(P^2)f^i} = \frac{\alpha(P^2)}{P^2 - \Pi_T^i|_{\alpha(P^2)}} \end{aligned}$$



- effective IR cut-off $\mu^2 \sim \Pi_T(0^-) \sim m_D^2$
- running coupling $\alpha(\mathbf{t})$ at virtuality

Running coupling

1-loop approximation $\alpha(t) = \frac{1}{4\pi\beta_0 \ln(-t/\Lambda^2)}$

- QED

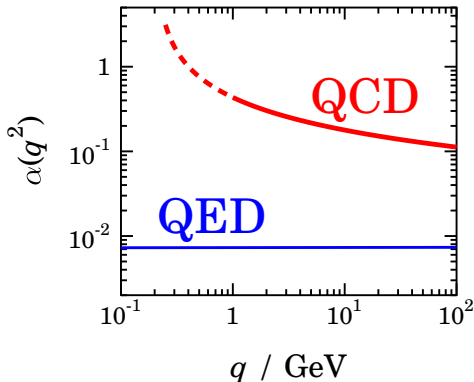
$$\beta_0 = -\frac{4}{3}/(4\pi)^2$$

$$\Lambda \sim 10^{xx} \text{ GeV}$$

- QCD

$$\beta_0 = (11 - \frac{2}{3}n_f)/(4\pi)^2$$

$$\Lambda \approx 0.2 \text{ GeV}$$



QCD: **running coupling crucial** at *probable* scales

- reconsider t -integration in energy loss

$$\begin{aligned} \Phi \int_{t_1}^{t_2} dt \frac{d\sigma_{js}}{dt} \Delta E &= -\frac{\pi C_{js}/k}{(4\pi\beta_0)^2} \int_{t_1}^{t_2} \frac{dt}{t \ln^2(|t|/\Lambda^2)} = \frac{\pi C_{js}/k}{(4\pi\beta_0)^2} \frac{1}{\ln(|t|/\Lambda^2)} \Big|_{t_1}^{t_2} \\ &= \frac{\pi C_{js}/k}{4\pi\beta_0} \left[\alpha(\mu^2) - \alpha(t_1) \right] \quad \text{instead of} \quad \frac{\pi C_{js}}{k} \alpha^2 \ln \frac{t_1}{\mu^2} \end{aligned}$$

for hard jets, due to $|t_1| \sim TE \gg \mu^2 \sim m_D^2 \sim \alpha T^2$, independent of E

- QCD collisional energy loss (with running coupling)

$$\frac{dE_j}{dx} = \pi \frac{\alpha(\mu^2)}{4\pi\beta_0} \sum_s C_{js} \int_{k^3} \frac{\rho_s(k)}{k} \rightarrow \left(\frac{2}{3}\right)^{\pm 1} 2\pi \left(1 + \frac{1}{6} n_f\right) T^2 \frac{\alpha(\mu^2)}{4\pi\beta_0}$$

- parametric enhancement: order α instead of α^2
- scale of running coupling is $\mu \sim m_D \sim \sqrt{\alpha} T$ instead of $\mathcal{O}(T)$
- independent of jet energy instead of $\sim \ln E$

Debye mass describes screening of static (color) electric fields

- defined as pole of longitudinal propagator

$$m_D^2 = \Pi'(0, p) \Big|_{p^2 = -m_D^2}$$

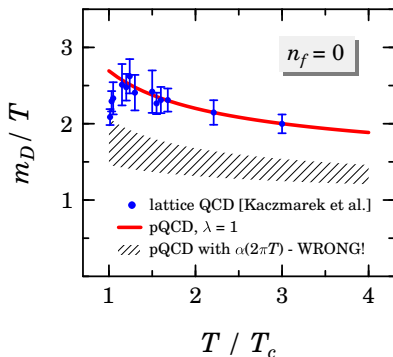
renormalize, express pole in terms of **running coupling**

$$P^2 = \# \alpha(P^2) f(p_0, p)$$

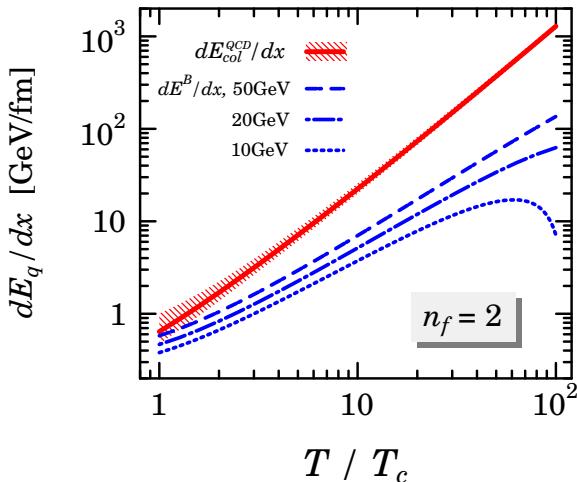
- implicit formula** for Debye mass

$$m_D^2 = \frac{1}{3} N_c \left(1 + \frac{1}{6} n_f\right) 4\pi \alpha(m_D^2) T^2$$

[AP, hep-ph/0601119]



$$\lambda_{fit}^{(0)} \approx \frac{1}{1.14} = \frac{\Lambda^{(0)}}{T_c^{(0)}} \Big|_{lattice}$$



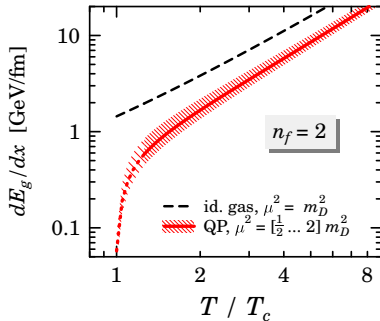
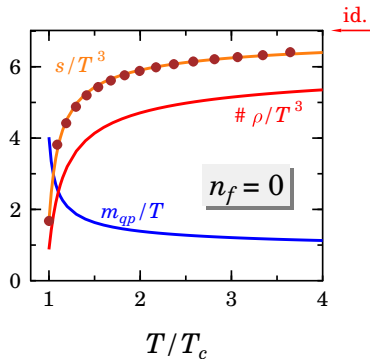
- $T_c = 160$ MeV
- $\Lambda = 205$ MeV
(consistent with lattice results for both $m_D(T)$ and $V(r)$ at $T = 0$)
- $\mu^2 = [\frac{1}{2}, 2] m_D^2$
- dE^B/dx calc. with $\alpha(2\pi T)$, and $\mu = m_D$

sizable enhancement, even near T_c , compared to previous estimates

Extrapolations, estimates

bulk properties of sQGP described by **quasiparticles** with mass $m_{qp} \sim \sqrt{\alpha} T$
 [AP, Kämpfer, Soff]

$$\frac{dE}{dx} \sim \int_{k^3} \rho(k) \Phi \int dt \frac{d\sigma}{dt} \Delta E$$



near T_c : large coupling \rightarrow reduced phase space density \rightarrow **small** dE/dx

- leading order QCD collisional energy loss

$$\frac{dE_{col}^{QCD}}{dx} \sim \alpha(m_D) T^2$$

- independent of jet energy
- non-Abelian enhancement compared to QED analog

- estimates for HIC's (based on parameters consistent with lattice QCD)
 - for $T \lesssim 1.5T_c$: strongly interacting, yet **transparent QGP**
⇒ no quenching at SPS
 - for $T \gtrsim 1.5T_c$: considerably **larger than previous estimates**
⇒ strong collisional quenching at RHIC/LHC